[1]E. H. Lieb and F. Y. Wu, Phys. Rev. Letters 20, 1445 (1968). Although there is some question of the completeness of the eigenfunctions of this paper, this solution does appear to give all of the finite-energy eigensolutions of Eq. (1) in the infinite U limit.

[2]A. B. Harris and R. V. Lange, Phys. Rev. 157, 292 (1967).

[3]H. S. Jarrett, W. H. Cloud, R. J. Bouchard, S. R. Butler, C. G. Frederick, and J. L. Gillson, Phys. Rev. Letters 21, 617 (1968); J. Appl. Phys. 40, 1258 (1969).

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ERRATA

Alpha-Particle Stopping Cross Section in Solids from 400 keV to 2 MeV, W. K. Chu and D. Powers [Phys. Rev. 187, 478 (1969)]. The authors are grateful to W. D. Mackintosh for kindly pointing out an error in the empirical formula given on p. 490, left column, second line from the bottom. The formula should read

$$\epsilon_\alpha = \epsilon_\alpha(E, Z_2) = 1.66(A'/E)\ln(B'E),$$

instead of the constant 1.66 being omitted. The correct form was used in the calculation for generating the curves appearing in Figs. 4–7.

Linear Chain Antiferromagnetism in $CsMnCl_3 \cdot 2H_2O$. T. Smith and S. A. Friedberg [Phys. Rev. 176, 660 (1968)]. The labels of the a and b axes have been inadvertently interchanged in Figs. 2, 3, and 4 of this paper and, consequently, also in the discussion of the results (Sec. IV) and in the abstract. This was kindly pointed out to us by Professor J. A. Cowen and verified independently in this laboratory by Dr. H. Kobayashi. As a result, the anisotropy axis above ~9 °K (taken as the z axis in the calculation on p. 663) is the a axis. The preferred axis of antiferromagnetic spin alignment in the three-dimensionally ordered phase ($T < 4.8$ °K) should be the b axis. Dr. Kobayashi's measurements also establish the anomalous susceptibility near 1 °K to be an impurity effect.

Equation (9) should read

$$\chi_n = \frac{Ng^2\mu_B{}^2 S(S+1)}{(n+1)3kT}\left\{(n+1)\frac{1+u}{1-u} - 2u\frac{1-u^{n+1}}{(1-u)^2}\right\}.$$

Magnetic Susceptibility of $FeCl_2 \cdot 4H_2O$ from 0.35 to 4.2 °K. J. T. Schriempf and S. A. Friedberg [Phys. Rev. 136, A518 (1964)]. To be consistent with the sign convention for D and E employed in Eq. (1) and elsewhere in this paper, Eqs. (14a)–(14c) should read $g_{x'} = g - (2D/3\lambda) + (2E/\lambda)$, $g_{y'} = g - (2D/3\lambda) - (2E/\lambda)$, $g_{z'} = g + (4D/3\lambda)$, respectively. The splitting factors given on p. A524 should thus read $g_{x'} = 2.18$, $g_{y'} = 2.17$, $g_{z'} = 2.22$.